A NOTE ON SINGULARITIES IN SEMILINEAR PROBLEMS

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ABSTRACT. We consider the equation $\Delta u - \frac{1}{2}x.\Delta u - \frac{u}{q-1} + u^q = 0$, for q > 1. We study the isolated singularities and present a nonlinear technique, and give a complete classification.

1. Introduction

In this note we study the isolated singularities of the positive solutions of the following nonlinear equation:

(1.1)
$$\Delta u - \frac{1}{2}x \cdot \nabla u - \frac{u}{q-1} + u^q = 0.$$

Let Ω be an open subset of \mathbb{R}^N , $N \geq 2$, containing 0, $\Omega' = \Omega \setminus \{0\}$ and q > 1. We are concerned with the following question:

If $u \in C^2(\Omega')$ is a positive solution of (1.1) in Ω' , what can be said about u(x) and about the equation as $|x| \to 0$?

It is well known [15] that (1.1) has no nontrivial globally bounded solution for N=1, 2 or N>2, $q\leq \frac{N+2}{N-2}$ except the constant solution $u=(\frac{1}{q-1})^{1/(q-1)}$. If we look for a specific solutions of (1.1) under the form

$$u(r) = \alpha r^{\beta},$$

then we get

$$\beta = -\frac{1}{q-1}$$
 and $\alpha = \lambda_{N,q} = \left\{ \frac{2}{q-1} \left(N - \frac{2q}{q-1} \right) \right\}^{1/(q-1)}$

where it is clear that $\lambda_{N,q}$ only exists when N>2 and $q>\frac{N}{N-2}$. It follows by substitution that $u(r)=\lambda_{N,q}r^{\beta}$ is also a solution of the Emden-Fowler equation

$$(1.2) \Delta w + w^q = 0.$$

This equation has been intensively studied. When N > 2 two critical values $\frac{N}{N-2}$ and $\frac{N+2}{N-2}$ appear. The first studies in the radial case are due to Emden, then Fowler [3, 5, 6], Brezis and Lions [2] and Lions [22] (see also [18, 1]).

Let us briefly describe our results. We have two cases.

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First case. $1 < q < \frac{N}{N-2}$. Then any positive solution of (1.1) in Ω' satisfies the equation $-\Delta u + \frac{1}{2}x \cdot \nabla u + \frac{u}{q-1} - u^q = \alpha \delta_0$ in $D'(\Omega)$ with $\alpha \ge 0$. Furthermore if $\alpha = 0$, then the "singularity is removable".

Second case. $\frac{N}{N-2} < q < \frac{N+2}{N-2}$. Then for any positive radial solution of (1.1) in Ω' we have $\lim_{r \to 0} r^{2/(q-1)} u(r) = l \in \{0, \lambda_{N,q}\}$ as $r \to 0$, where r = |x|. If l = 0 and q > 2, the "singularity is removable". Note that if $q \ge \frac{N}{N-2}$, any solution of (1.1) satisfies

$$-\Delta u + \frac{1}{2}x \cdot \nabla u + \frac{u}{q-1} - u^q = 0, \quad \text{in } D'(\Omega).$$

2. Main results

Let N > 2 and set $\Omega = B_R(0) = \{x \in \mathbb{R}^N, |x| < R\}, R > 0$ and $\Omega' = \Omega \setminus \{0\}$. Our main result is the following:

Theorem 2.1. Assume that $u \in L^1_{Loc}(\Omega')$,

(2.1)
$$\Delta u - \frac{1}{2}x.\nabla u \in L^1_{Loc}(\Omega'), \quad in \ D'(\Omega'),$$

(2.2)
$$u \ge 0, \qquad \Delta u - \frac{1}{2} x. \nabla u \le au + f \text{ a.e. in } \Omega,$$

where a is a nonnegative constant and $f \in L^1_{Loc}(\Omega)$. Then

$$(2.3) u \in L^1_{\text{Loc}}(\Omega),$$

and there exists $h \in L^1_{loc}(\Omega)$, and $\alpha \ge 0$ such that

(2.4)
$$-\Delta u + \frac{1}{2}x \cdot \nabla u = h + \alpha \delta_0 \quad \text{in } D'(\Omega) \, .$$

For the proof, we use a nonlinear technique introduced by Serrin [23], [24] and a linear method by H. Brezis and P. L. Lions [2], [22].

The proof is divided into five steps.

Step 1. We claim that

$$x.\nabla u = r\frac{\partial u}{\partial r}$$

where $x = (r, \sigma), r = |x|$.

Proof of the claim. Writing

$$x.\nabla u = \sum_{i=1}^{N} x_i \left(\frac{\partial u}{\partial r} \frac{\partial r}{\partial x_i} + \sum_{j=1}^{N-1} \frac{\partial u}{\partial \sigma_j} \frac{\partial \sigma_j}{\partial x_i} \right),$$
$$x.\nabla u = r \frac{\partial u(r, \sigma)}{\partial r} + \sum_{j=1}^{N-1} \frac{\partial u}{\partial \sigma_j} \sum_{i=1}^{N} x_i \frac{\partial \sigma_j}{\partial x_i},$$

as in [17], we have

$$\frac{\partial \sigma_{j}}{\partial x_{i}} = \begin{cases} \frac{x_{i}x_{j+1}}{r_{j}r_{j+1}^{2}} & \text{if } i \leq j, \\ -\frac{r_{j}}{r_{j+1}^{2}} & \text{if } i = j+1, \\ 0 & \text{if } i > j+1, \end{cases}$$

where $r_j^2 = x_1^2 + x_2^2 + \dots + x_j^2$, $r_N^2 = r^2$. Thus

$$\sum_{i=1}^{N} x_i \frac{\partial \sigma_j}{\partial x_i} = 0.$$

And then

$$(2.5) x.\nabla u = r\frac{\partial u}{\partial r}.$$

Step 2. $u \in L^1_{\mathrm{loc}}(\Omega)$. We begin by considering the average

$$\overline{u}(r) = \frac{1}{\omega_{N-1}} \int_{S^{N-1}} u(r, \sigma) d\sigma, \qquad 0 < r < R,$$

where ω_{N-1} is the volume of the sphere S^{N-1} . It follows from (2.1), (2.2) and (2.5) that

(2.6)
$$\Delta \overline{u} - \frac{1}{2} r \overline{u}_r \in L^1_{Loc}(0, R), \qquad \overline{u} \geq 0,$$

(2.7)
$$\Delta \overline{u} - \frac{1}{2} r \overline{u}_r \le a \overline{u} + \overline{f} \quad \text{on } (0, R),$$

and

(2.8)
$$\frac{1}{r^{N-1}}(r^{N-1}K(r)\overline{u}_r)_r \le a\overline{u} + \overline{f}, \quad \text{for } r \in (0, R),$$

where $K(r) = e^{-r^2/4}$. In particular $\overline{u} \in C^1(0, R)$.

Let R' < R be fixed. Integrating (2.8) over (r, R') we find as in [2] that

$$(2.9) u \in L^1_{loc}(\Omega)$$

and

$$(2.10) \overline{u}(r) \le \frac{C}{r^{N-2}} + C.$$

Step 3. Let $g(x) = -\operatorname{div}(K(x)\nabla u)$ a.e. in Ω , where $K(x) = e^{-|x|^2/4}$. Then $g \in L^1_{loc}(\Omega)$ and for any $\varphi \in D(\Omega)$, $0 \le \varphi \le 1$, $\varphi \equiv 1$ near x = 0,

(2.11)
$$\int_{\Omega} g \varphi^2 dx \le -\int_{\Omega} u. \operatorname{div}(K(x) \nabla \varphi^2) dx.$$

Proof of the step 3. From the definition of g we have

(2.12)
$$\int_{\Omega} g \psi \, dx = \int_{\Omega} K(x) \nabla u \nabla \psi \, dx,$$

for any $\psi \in W^{1,\infty}(\Omega')$ with compact support in Ω' . Set

$$P_k(t) = \begin{cases} 1 & \text{if } t > 1 + k, \\ t - k & \text{if } k \le t \le 1 + k, \\ 0 & \text{if } t < k \end{cases}$$

for any $t \ge 0$ and $k \ge 0$.

Let $0<\rho< R$ and $\varepsilon<\frac{\rho}{2}$. Let $\varphi\in D(\Omega)$ be such that $0\leq \varphi\leq 1$, $\varphi\equiv 1$ on $B_{\rho}=\{x\in \mathbf{R}^{N}\,,\,|x|<\rho\}$ and $\psi_{\varepsilon}=\eta_{\varepsilon}\varphi$ with $\eta_{\varepsilon}\in C^{\infty}(\Omega)\,,\,0\leq\eta_{\varepsilon}\leq 1\,,\,\eta_{\varepsilon}\equiv 0$ if $|x|<\varepsilon$, $\eta_{\varepsilon}(x)\equiv 1$ if $|x|>2\varepsilon$ and $|\nabla\eta_{\varepsilon}|\leq \frac{\varepsilon}{\varepsilon}$. We have, from (2.12),

(2.13)
$$\int_{\Omega} g(1 - P_k(u)) \psi_{\varepsilon}^2 dx + \int_{\{k < u < 1 + k\}} K(x) |\nabla u|^2 \psi_{\varepsilon}^2 dx$$
$$= \int_{\Omega} (1 - P_k(u)) K \nabla u \nabla \psi_{\varepsilon}^2 dx.$$

Since

$$\sum_{k=0}^{n} (1 - P_k(t))^+ = (n+1-t)^+,$$

we have

$$\begin{split} &\int_{\{u < n+1\}} g(n+1-u)^+ \psi_{\varepsilon}^2 \, dx + \int_{\{u < 1+n\}} K(x) |\nabla u|^2 \psi_{\varepsilon}^2 \, dx \\ &\leq \int_{\{u < n+1\}} (n+1-u) K \nabla u \nabla \varphi^2 \eta_{\varepsilon}^2 \, dx \\ &\quad + 2 \int_{\{u < 1+n\} \cap B_{2\varepsilon}} (n+1-u) K(x) |\nabla u| \psi_{\varepsilon} |\nabla \eta_{\varepsilon}| \, dx \, . \end{split}$$

Now for any real h>0, we have $n+1-u(x)>n+1\frac{h}{h+1}$ a.e. in $\{u<\frac{n+1}{h+1}\}$. Thus, dividing by n+1 and using the Hölder inequality, for any $\beta>0$,

$$\frac{h}{h+1} \int_{\{u < \frac{n+1}{h+1}\}} g \psi_{\varepsilon}^{2} + \frac{1}{n+1} \int_{\{u < 1+n\}} K |\nabla u|^{2} \psi_{\varepsilon}^{2}
\leq \int_{\{u < 1+n\}} \left(1 - \frac{u}{n+1}\right) K \nabla u \nabla (\varphi^{2}) \eta_{\varepsilon}^{2}
+ \beta^{2} \int_{\{u < 1+n\}} K(x) |\nabla u|^{2} \psi_{\varepsilon}^{2} + \beta^{-2} \int_{B_{2\varepsilon}} K(x) |\nabla \eta_{\varepsilon}|^{2}.$$

Let $\beta^2 = \frac{1}{2(n+1)}$; then

$$\frac{h}{h+1} \int_{\{u < \frac{n+1}{h+1}\}} g \psi_{\varepsilon}^{2} + \frac{1}{2(n+1)} \int_{\{u < 1+n\}} K |\nabla u|^{2} \psi_{\varepsilon}^{2} \\
\leq \int_{\{u < 1+n\}} \left(1 - \frac{u}{n+1}\right) K \nabla u \nabla (\varphi^{2}) \eta_{\varepsilon}^{2} + 2(n+1) C \varepsilon^{N-2}.$$

From Fatou's Lemma—which is valid since $g \ge -au - f \in L^2_{loc}(\Omega)$ —we deduce as $\varepsilon \to 0$, $n \to +\infty$, $h \to +\infty$ that $g\varphi^2 \in L^1(\Omega)$ and satisfies (2.11).

Step 4. Now since $u \in L^1(\Omega)$, we can define the distribution

$$T = -\operatorname{div}(K\nabla u) - g$$
 in $D'(\Omega)$.

Then as in [2], we have

$$T = \sum_{|p| \le m} c_p D^p \delta_0.$$

Let $\psi \in D(B_R)$ be any fixed function such that

$$(-1)^{|p|}D^p\psi(0)=c_p$$
 for every $|p|\leq m$

and

$$\psi_{\varepsilon}(x) = \psi\left(\frac{x}{\varepsilon}\right)$$
.

Then

(2.14)
$$\int_{B_R} u \operatorname{div}(K \nabla \psi_{\varepsilon}) = \int_{B_R} g \psi_{\varepsilon} + \sum_{|p| \le m} \frac{c_p^2}{\varepsilon^{|p|}}.$$

On the other hand we have

$$\left| \int_{B_{\mathbb{R}}} u \operatorname{div}(K \nabla \psi_{\varepsilon}) \right| \leq \frac{C}{\varepsilon^2} \int_{B_{\mathbb{R}}} u \, dx + CR \int_{B_{\mathbb{R}}} u \, dx$$

and therefore

$$\left| \int_{B_R} u \operatorname{div}(K \nabla \psi_{\varepsilon}) \right| \leq \frac{C}{\varepsilon^2} \int_0^{R\varepsilon} \overline{u} r^{n-1} dr + CR \int_0^{R\varepsilon} \overline{u} r^{n-1} dr.$$

We deduce from (2.10) that $|\int_{B_R} u \operatorname{div}(K \nabla \psi_{\varepsilon})| \leq C$. Comparing this with (2.14) we conclude that $c_p = 0$ when $|p| \geq 1$.

Finally we choose $\eta \in D(\Omega)$, $0 \le \eta \le 1$, and $\eta \equiv 1$ near x = 0.

We have

$$\langle T, \eta^2 \rangle = c_0 = \int_{\Omega} u \operatorname{div}(K \nabla \eta^2) - \int_{\Omega} g \eta^2,$$

and hence $c_0 \ge 0$ from (2.11).

Step 5. The end of the proof of Theorem 2.1.

Let $h(x) = g(x)e^{|x|/4}$ a.e. in Ω ; then $h \in L^1_{loc}(\Omega)$ and $-\Delta u + \frac{1}{2}x \cdot \nabla u = h + c_0 \delta_0$ in $D'(\Omega)$.

3. The subcritical case

Now we return to equation (1.1). Let μ be the fundamental harmonic function in $\mathbb{R}^N \setminus \{0\}$, N > 2, that is

(3.1)
$$\mu(x) = \frac{1}{N(N-2)\omega_N} |x|^{2-N}.$$

Theorem 3.1. Let N>2, $1< q<\frac{N}{N-2}$. Let $u\in C^2(\Omega')$ be a nonnegative solution of equation

(3.2)
$$\Delta u - \frac{1}{2}x \cdot \nabla u - \frac{u}{q-1} + u^q = 0 \quad in \ \Omega' = \Omega \setminus \{0\}.$$

Then

- (i) either u can be extended to a smooth solution of (3.2) in Ω , or
- (ii) there exists $\alpha > 0$ such that $\lim_{x\to 0} \frac{u(x)}{\mu(x)} = \alpha$ and u satisfies the equation

$$(3.3) -\Delta u + \frac{1}{2}x \cdot \nabla u + \frac{u}{q-1} - u^q = \alpha \delta_0 \quad in \ D'(\Omega).$$

Proof. From (3.2) we have $\Delta u - \frac{1}{2}x \cdot \nabla u \leq \frac{u}{q-1}$; hence from Theorem 2.1 we have

$$-\Delta u + \frac{1}{2}x \cdot \nabla u + \frac{u}{q-1} - u^q = \beta \delta_0, \qquad u \in L^1_{loc}(\Omega), \ u^q \in L^1_{loc}(\Omega).$$

Moreover equation (3.2) can be written in the form

(3.4)
$$\operatorname{div}(K\nabla u) + d(x)u = 0 \quad \text{in } \Omega'$$

where $d(x)=k(x)(u^{q-1}-\frac{1}{q-1})$. Since $q<\frac{N}{N-2}$ we can find an $\varepsilon>0$ such that $d\in L^{N/(2-\varepsilon)}_{loc}(\Omega)$. We then deduce from [24] that if u is singular at 0 there exists a constant C>0 such that

$$(3.5) C\mu \le u \le \frac{1}{C}\mu \quad \text{near } 0.$$

As in Guedda-Veron [18], we prove by scaling that there is an $\alpha > 0$ such that

$$\lim_{x \to 0} \frac{u(x)}{\mu(x)} = \alpha;$$

we get that

(3.6)
$$\lim_{x \to 0} |x|^{N-1} \nabla u(x) = \frac{\alpha}{N\omega_N} \xi, \quad \text{where } \xi = \lim_{x \to 0} \frac{x}{|x|}$$

and then (3.3).

4. THE SUPER CRITICAL CASE

We still assume that $\Omega = B_R(0) = \{x \in \mathbf{R}^N, |x| < R\}$ and $\Omega' = \Omega \setminus \{0\}$. In this section we present some results concerning the isolated singularities of the positive solutions of (1.1), that is, of

$$-\Delta u + \frac{1}{2}x.\nabla u + \frac{u}{a-1} - u^q = 0 \quad \text{in } \Omega',$$

where

$$\frac{N}{N-2} < q < \frac{N+2}{N-2}$$
 and $N > 2$.

If we look for the solution of (3.2) under the form $u(x) = u(|x|) = \alpha |x|^{\beta}$, then we get

$$\beta = -\frac{1}{q-1}$$
 and $\alpha = \lambda_{N,q} = \left(\frac{2}{q-1}\left(N - \frac{2q}{q-1}\right)\right)^{1/(q-1)}$.

Theorem 4.1. Let N > 2, $q \ge \frac{N}{N-2}$. Let $u \in C^2(\Omega')$ be a nonnegative solution of (4.1) in $D'(\Omega')$. Then we have $\alpha = 0$ in (3.3), i.e.

$$-\Delta u + \frac{1}{2}x \cdot \nabla u + \frac{u}{q-1} - u^q = 0 \quad in \ D'(\Omega).$$

Theorem 4.2. Under the assumptions of Theorem 4.1, if

(4.2)
$$u^{(q-1)N/2} \in L^1_{loc}(\Omega)$$
,

then u can be extended to Ω as a solution of (4.1) in Ω .

The proofs of Theorems 4.1 and 4.2 are the same as in Guedda and Veron [18] and they are omitted.

The main result of this section is the following

Theorem 4.3. Let N > 2 and $\frac{N}{N-2} < q < \frac{N+2}{N-2}$ and u be a positive radial solution of (4.1).

Then we have

- (i) $r^{2/(q-1)}u(r)$ has a limit l as $r \to 0$ and $l \in \{0, \lambda_{N,q}\}$.
- (ii) If l = 0 and q > 2, then u can be extended to Ω as a $C^2(\Omega)$ solution of (4.1) in Ω .

We need the following

Lemma 4.1. Let u be a positive radial solution of (4.1) in Ω' . Then $r^{2/(q-1)}u(r)$ is bounded as $r \to 0$.

Proof. Let

(4.5)
$$t = \log(r^{\gamma}), \quad v(t) = r^{2/(q-1)}u(r), \quad \gamma = \frac{4}{q-1} + 2 - N > 0;$$

then (4.1) takes the equivalent form

(4.6)
$$\gamma^{2}(v_{tt} + v_{t}) + \frac{1}{2}\gamma e^{-2t/\gamma}v_{t} + \lambda v + v^{q} = 0$$

where $\lambda = \frac{2}{q-1}(\frac{2}{q-1} + 2 - N) < 0$. Let

$$L(t) = \frac{1}{2}\gamma^2 v_t^2 + \lambda \frac{v^2}{2} + \frac{v^{q+1}}{q+1}, \quad \text{for } t \in [0, +\infty[.$$

From (4.6) the function L is nonincreasing. Hence for any $t_0 \le t < +\infty$

$$(4.7) \frac{1}{2} \gamma^2 v_t^2 + \lambda \frac{v^2}{2} + \frac{v^{q+1}}{q+1} \le L(t_0).$$

Hence v is bounded, as q > 1.

Lemma 4.2. Suppose $\frac{N}{N-2} < q < \frac{N+2}{N-2}$. Let v be any solution of (4.6). Then $\lim_{t\to +\infty} v_t = 0$ and v has a limit l at $+\infty$ and

$$(4.8) l(\lambda + l^{q-1}) = 0.$$

Proof. Using (4.6), (4.7) and Lemma 4.1 we have

$$(4.9) v, v_t, v_{tt} \text{ and } v_{ttt} \text{ are bounded.}$$

Hence $v_t \in L^2(t_0, +\infty)$. This implies $v_{tt} \in L^2(t_0, +\infty)$. It follows from this and (4.9) that

$$\lim_{t\to+\infty}v_t=\lim_{t\to+\infty}v_{tt}=0.$$

For any $t \ge \tau \ge t_0$, we have

$$|v(t)-v(\tau)| \leq ||v_t||_{L^2(t_0,+\infty)} \cdot \sqrt{(t-\tau)},$$

and hence v has a limit l at $+\infty$. From (4.6) we get (4.8).

Proof of Theorem 4.3. From Lemmas 4.1 and 4.2 $r^{2/(q-1)}u(r)$ has a limit l where l=0 or $\lambda_{N,q}$.

Let v be as in (4.5). Assume that l=0. The proof of (ii) is divided into two steps.

Step 1. Assume v decreases to 0 as t tends to $+\infty$; then

$$(4.10) v(t) \le Ce^{-t},$$

for $t \ge 0$ and some C > 0. In fact, from (4.6) we have

(4.11)
$$\gamma^{2}(v_{t}+v) = \int_{t}^{+\infty} \frac{1}{2} \gamma e^{-2t/\gamma} v_{s} + \lambda v + v^{q} ds.$$

Since v and v_t are bounded, $\omega(t) = \frac{1}{2}\gamma e^{-2t/\gamma}v_s + \lambda v + v^q$ is integrable; then $\lambda v + v^q$ is integrable and then v is integrable. Since $v_t \le 0$ and $\lambda < 0$, we have

$$\gamma^2(v_t+v) \le \int_t^{+\infty} v^q \, ds \le v^{q-1} \int_t^{+\infty} v \, ds;$$

then

$$(4.12) v_t + v \le Cv^{q-1}, for t large,$$

which implies, if q > 2, (4.10). Moreover, Theorem 4.2 implies that if v satisfies (4.10), then u can be extended to Ω as a regular solution of (1.1) in Ω .

Step 2. Assume that v is not asymptotically monotone. Then there exists a sequence $\{t_n\}$ such that $v_t(t_n) = 0$, $v(t_{2n})$ is local minimum and $v(t_{2n+1})$ is local maximum, and (from the equation)

$$0 < v(t_{2n}) < \lambda_{N,q} < v(t_{2n+1})$$

which is not possible since $\lim_{t\to +\infty} v(t) = 0$, which ends the proof.

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